

### 3. Magneto statics & Amphere's law

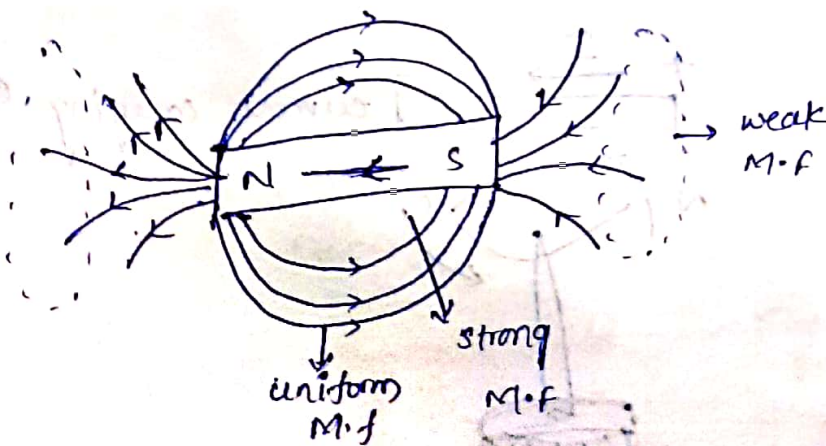
#### Introduction:-

In this chapter, the magnetic field is introduced the relation of steady magnetic fields to its source is more complicated when compared to the relation b/w electrostatic field to its source.

The development of motors, transformers, high speed vehicles, Television etc. involve magnetic phenomena and play an important role in every day life.

There is a great similarity in the equations derived in electric fields and Magnetic fields.

#### \* Magnetic field Intensity:-



- The region surrounding a Magnet in which the force of the Magnet can be experienced a field that is called as Magnetic field. it is indicated by  $H$ .  $\text{AmpTurn/m}$

## Note:-

- The lines flow from North - South (outside the magnet) and South - North (inside the magnet)
- The lines show both direction and strength.

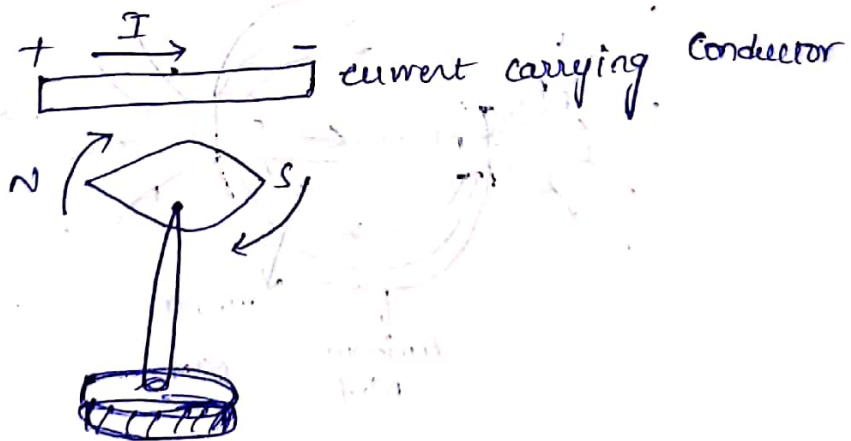
## Magnetic flux density (B):-

It is defined as the magnetic flux per unit surface area. It is a vector quantity.

$B$  units  $\text{wb/m}^2$  or Tesla

Mathematically,  $B = \frac{\phi}{A}$   $\text{wb/m}^2$  (or) Tesla

## \*Oersted's Experiment:-



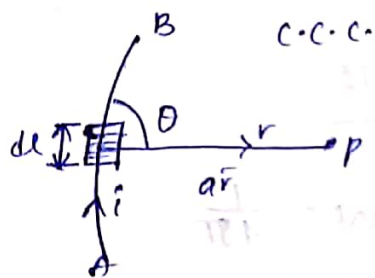
- Oersted conducted an experiment in which a current carrying conductor was taken. A compass needle was kept under this conductor as shown in figure.

- when there was no current through the conductor, the needle was pointing along north-south of the earth.

- But when the conductor carries current then the needle was attracted to the conductor or repulsion to the conductor, it moved and tended to stand at right angles to the conductor.

- From this experiment Oersted shows that an electric current produces a magnetic field.

\* Biot-Savart's law:-



- 1820's Biot and Savart conducted a series of experiments on steady magnetic field. finally they conclude that the required point  $\vec{B}$  is obtained from the following relation  
[ $\vec{H} \text{ or } \vec{B}$ ]

i) The differential magnetic field intensity is directly proportional to current carrying through the conductor.

$$dH \propto I$$

ii) The differential MFI is directly proportional to differential length

$$dH \propto dL$$

iii) The differential MFI is directly proportional to sine angle

$$dH \propto \sin \theta$$

iv) The differential MFI is Inversely proportional to square of distance from differential length to point P.

$$dH \propto 1/r^2$$

$$\therefore dH \propto \frac{I dl \sin \theta}{r^2}$$

$$dH = k \frac{I dl \sin \theta}{r^2}$$

where  $k = \text{constant} = \frac{1}{4\pi}$

$$\Rightarrow dH = \frac{I dl \sin \theta}{4\pi r^2} \bar{a}_r$$

If  $\theta = 90^\circ$

$$dH = \frac{I dl}{4\pi r^2} \bar{a}_r$$

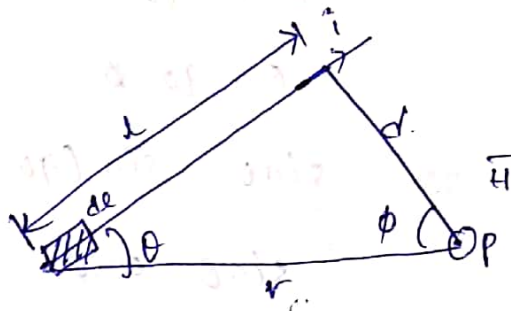
$$H = \int \frac{I dl}{4\pi r^2} \sin \theta \bar{a}_r$$

we know that  $B = \mu H$

$$B = \mu \int \frac{I dl}{4\pi r^2} \sin \theta \bar{a}_r$$

$$B = \frac{\mu}{4\pi} \int \frac{I dl \sin\theta}{r^2} \vec{a}_r$$

\* Magnetic field intensity due to current carrying conductor or wire:-



consider a straight wire of length 'l' carrying a steady current \$I\$. here we find the magnetic field intensity \$\vec{H}\$ at point 'P', with a distance of 'd' from the wire as shows in fig:

from the figure;

$$\sin\theta = d/r = \cos\phi$$

$$\cos\theta = l/r = \sin\phi$$

$$\tan\phi = \frac{l/r}{d/r} = l/d$$

$$l = d \tan\phi$$

diff. w.r.t. \$\phi\$

$$\frac{dl}{d\phi} = d \sec^2\phi$$

$$dl = d \sec^2\phi d\phi$$

$$r^2 = l^2 + d^2$$

$$r = \sqrt{l^2 + d^2}$$

$$\phi + \theta = 90^\circ$$

$$\phi + \theta = 90^\circ$$

$$\phi = 90 - \theta$$

$$\theta = 90 - \phi$$

$$\sin \phi = \sin(90 - \theta)$$

$$\sin \theta = \sin(90 - \phi)$$

$$\sin \phi = \cos \theta$$

$$\sin \theta = \cos \phi$$

w.k.t

The MFI at any point 'p' is,

$$d\vec{H} = \frac{I dl \sin \theta}{4\pi r^2} \vec{a}_r$$

$$d\vec{H} = \frac{I dl \cos \phi}{4\pi r^2} \vec{a}_r$$

$$\vec{H} = \int \frac{I dl \cos \phi}{4\pi r^2} \vec{a}_r$$

$$= \frac{I}{4\pi} \int \frac{dl \cos \phi}{r^2} \vec{a}_r$$

$$= \frac{I}{4\pi} \int \frac{d \sec^2 \phi d\phi \cos \phi}{(l^2 + d^2)} \vec{a}_r$$

$$= \frac{I}{4\pi} \int \frac{d \sec^2 \phi d\phi (d/r)}{l^2 + d^2} \vec{a}_r$$

$$= \frac{I d^2}{4\pi} \int \frac{\sec^2 \phi d\phi}{(l^2 + d^2)^{3/2}} \vec{a}_r$$

$$\bar{H} = \frac{2d^2}{4\pi} \int \frac{\sec^2 \phi d\phi}{(r^2 + d^2)^{3/2}} \cdot \bar{ar}$$

(dr)<sup>1/2</sup>  
(r^2)<sup>1/2</sup> · (r^2 + d^2)<sup>1/2</sup>

$$= \frac{2d^2}{4\pi} \int \frac{\sec^2 \phi d\phi}{d^3 \left[1 + \frac{r^2}{d^2}\right]^{3/2}} \bar{ar}$$

$$= \frac{2d}{4\pi d} \int \frac{\sec^2 \phi d\phi}{(1 + \tan^2 \phi)^{3/2}} \bar{ar}$$

$$= \frac{2}{4\pi d} \int \frac{\sec^2 \phi d\phi}{(\sec^2 \phi)^{3/2}} \bar{ar}$$

$$= \frac{2}{4\pi d} \int \frac{1}{\sec \phi} d\phi \bar{ar}$$

$$= \frac{2}{4\pi d} \int \cos \phi d\phi \bar{ar}$$

$$\bar{H} = \frac{2}{4\pi d} \sin \phi \bar{ar}$$

$$\bar{H} = \frac{2 \cos \theta}{4\pi d} \bar{ar}$$

Note:-

for infinite length i.e.,  $-\frac{\pi}{2}$  to  $+\frac{\pi}{2}$  the limits of integration. ∴ the MFI due to CCC is (or)  $\frac{2 \cos \theta}{4\pi d}$

$$\bar{H} = \frac{2 \sin \phi}{4\pi d} \bar{ar} = \frac{2(\sin \phi_2 - \sin \phi_1)}{4\pi d} \bar{ar}$$

$$= \frac{2}{4\pi d} \left[ \sin \frac{\pi}{2} + \sin \frac{\pi}{2} \right] = \frac{2}{4\pi d} = \frac{2 \cos \theta_1 - \cos \theta_2}{4\pi d}$$

$$\bar{H} = \frac{2}{2\pi d}$$

\* Magnetic permeability: - ( $\mu$ )

It is the ratio flux density ( $B$ ) and produced in the material or produced M.F.S.

$$\therefore B \propto H$$

$$B = \mu H$$

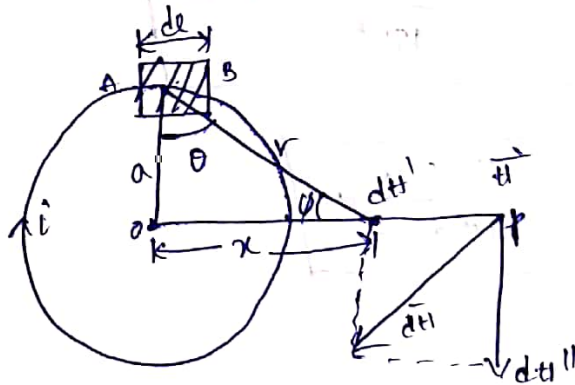
$$\boxed{\mu = B/H}$$

where  $\mu = \mu_0 \mu_r$

$\mu_0$  = Absolute permeability =  $4\pi \times 10^{-7} \text{ T/m}$

$\mu_r$  = relative permeability

\* Magnetic f.S due to circular coil: -



consider a circular loop of a radius  $a$  carrying a current  $i$  amp. we have to find the magnetic field intensity at point P on the axis of loop with a distance of  $x$  mt. from its centre  $O$ .

consider a segment AB whose elemental length is ' $dl$ ',

let  $r$  be the distance of the element from point P or  $\theta$  be



the angle b/w the direction element with line joining the elementary length dl.

- from Biot-Savart's law

$$\text{The MFI at point p is } d\vec{H} = \frac{\mu_0 I dl \sin\theta}{4\pi r^2} \vec{a}_r \rightarrow (1)$$

from the figure

$$r^2 = x^2 + a^2$$

$$r = \sqrt{x^2 + a^2}$$

$$\sin\theta = \frac{a}{r}$$

If  $\theta = 90^\circ$  then the equation (1) is,

$$d\vec{H} = \frac{\mu_0 I dl}{4\pi r^2} \vec{a}_r \rightarrow (2)$$

from the figure, the resultant vector  $d\vec{H}$

$$d\vec{H} = d\vec{H}' + d\vec{H}''$$

$d\vec{H}'' = 0$  because it is vertical component.

$\therefore$  The resultant vector  $d\vec{H} = d\vec{H}' \rightarrow (3)$

$$\text{w.k.t } d\vec{H} = d\vec{H} \sin\theta \rightarrow (4)$$

substitute (4) in (3)

$$d\vec{H} = d\vec{H} \sin\theta$$

$$d\vec{H} = \frac{\mu_0 I dl}{4\pi (x^2 + a^2)} \sin\theta \vec{a}_r$$

$$d\vec{H} = \frac{\mu_0 I dl}{4\pi (x^2 + a^2)} \cdot \frac{a}{r} \vec{a}_r$$

$$d\vec{H} = \frac{\mu_0 I a dl}{4\pi (x^2 + a^2)^{3/2}} \vec{a}_r$$

$$= \frac{\mu_0 I a dl}{4\pi (x^2 + a^2)^{3/2}} \vec{a}_r$$

$$\vec{H} = \int \frac{\mu_0 I a dl}{4\pi (x^2 + a^2)^{3/2}} \vec{a}_r$$

$$= \frac{\mu_0 I}{4\pi} \int \frac{dl \cdot a}{(x^2 + a^2)^{3/2}} \vec{a}_r$$

$$= \frac{\mu_0 I \cdot a}{4\pi (x^2 + a^2)^{3/2}} \int dl \cdot \vec{a}_r$$

$$= \frac{\mu_0 I \cdot a \cdot l}{4\pi (x^2 + a^2)^{3/2}} \vec{a}_r$$

$$\vec{H} = \frac{\mu_0 I \cdot a \cdot (2\pi a)}{4\pi (x^2 + a^2)^{3/2}} \vec{a}_r$$

$$\vec{H} = \frac{\mu_0 I a^2}{2 (x^2 + a^2)^{3/2}} \vec{a}_r$$

If the required point is at centre

$$\therefore x = 0$$

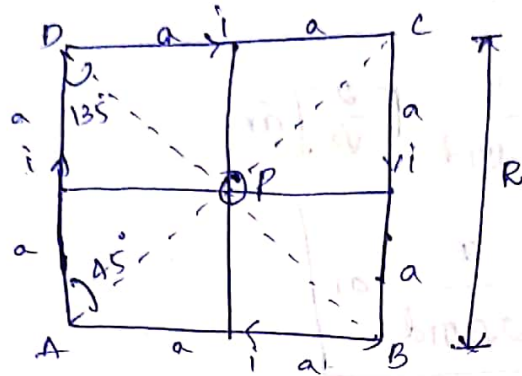
$$\vec{H} = \frac{\mu_0 I a^2}{2 a^3} \vec{a}_r$$

$$\vec{H} = \frac{\mu_0 I}{2a} \vec{a}_r$$

for  $n$  turns coils, the magnetic field intensity is

$$\vec{H} = \frac{IN}{2a} \vec{ar}$$

\* Magnetic field intensity at the centre of conductor or in the form of square:-



- let the length of the each side by  $R$  mt. and  $i$  denotes a steady direct current carried by loop. it is requires to determine the magnetic field intensity at point  $p$  at centre of square.

- consider the side 'AD' of the loop the field intensity at a point distance one form a current carrying conductor of finite length is given by,

$$\vec{H} = \frac{I d \cos \theta}{4\pi r d} \vec{ar} \rightarrow (1)$$

$$= \frac{I}{4\pi d} (\cos \alpha_1 - \cos \alpha_2) \vec{ar}$$

The side  $AB$  of the field intensity  $H_1$  at point  $P$

is given by

$$H_1 = \frac{I}{4\pi d} [\cos 45^\circ - \cos 135^\circ] \bar{a}_r$$

$$\alpha_1 = 45^\circ, \alpha_2 = 135^\circ$$

$$= \frac{I}{4\pi d} \left[ \frac{1}{\sqrt{2}} - \left[ -\frac{1}{\sqrt{2}} \right] \right] \bar{a}_r$$

$$= \frac{I}{4\pi d} \left[ \frac{2}{\sqrt{2}} \right] \bar{a}_r$$

$$\boxed{\bar{H}_1 = \frac{I}{2\sqrt{2}\pi d} \bar{a}_r}$$

it is clear that the field intensity  $H_1, H_2, H_3$  and  $H_4$  due to sides of AD, DC, CB, BA (due to direction) are equal.

The total field intensity at point P due to square loop ABCD is

$$\bar{H} = H_1 + H_2 + H_3 + H_4$$

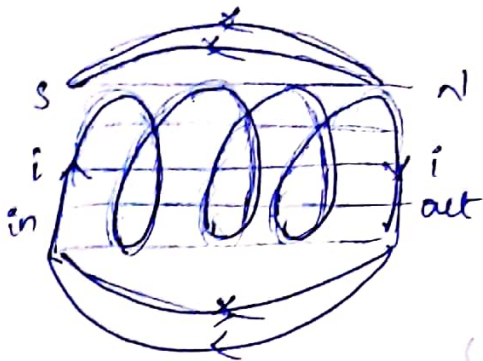
$$\Rightarrow \bar{H} = 4H_1$$

$$= 4 \left[ \frac{I}{2\sqrt{2}\pi d} \bar{a}_r \right]$$

$$\boxed{\bar{H} = \frac{\sqrt{2}I}{\pi d} \bar{a}_r}$$

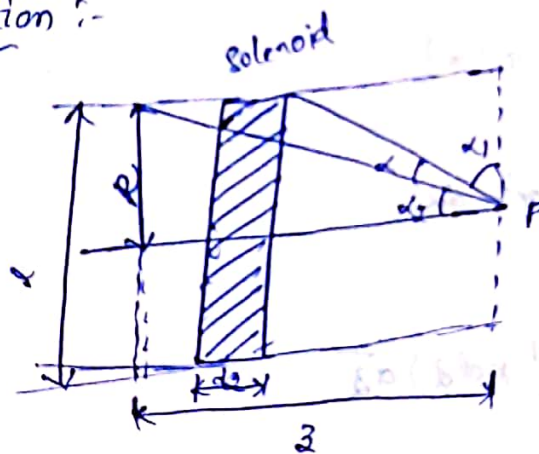
\* MFI due to solenoid:-

Solenoid:-



- A solenoid consists of long conducting wire made up of many loops packed closely together. When the current is passed through the solenoid then the magnetic field is produced.

Derivation:-



- consider a solenoid of  $n$ -turns of length  $l$ . Let  $R$  be the radius of solenoid.

- Let  $P$  be the point on the axis of the solenoid ( $z$ -axis)

- consider an element ' $dz$ ' at a distance ' $z$ ' from point  $P$ .

- No. of turns per unit length =  $N/l$ .

No. of turns in solenoid ( $dz$ ) =  $N/l \cdot dz \rightarrow \infty$

Magnetic field intensity for one turn at point P is,

$$\vec{H} = \frac{Ia^2}{2(x^2+a^2)^{3/2}} \cdot \vec{a}_r$$

in solenoid  $\vec{H} = \frac{IR^2}{2(z^2+R^2)^{3/2}} \cdot \vec{a}_z$

$$d\vec{H} = \frac{IR^2}{2(z^2+R^2)^{3/2}} \vec{a}_z \rightarrow (2)$$

from the fig;

$$\cot \alpha = z/R$$

diff w.r.t.  $\alpha$

$$\frac{dz}{d\alpha} = R(-\operatorname{cosec}^2 \alpha)$$

$$dz = -R \operatorname{cosec}^2 \alpha d\alpha$$

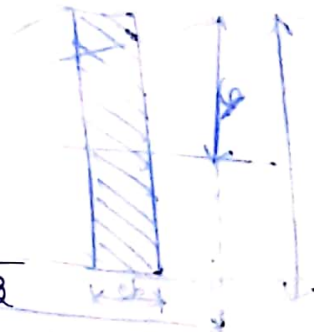
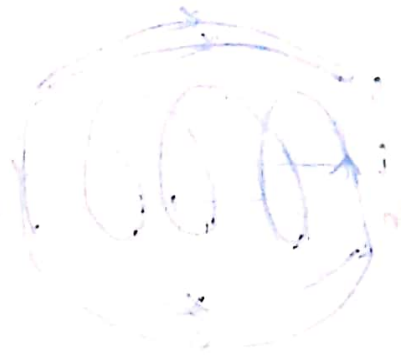
$$d\vec{H} = \frac{IR^2}{2(z^2+R^2)^{3/2}} \left( \frac{N}{l} \times dz \right) \vec{a}_z$$

$$= \frac{IR^2}{2(R^2+z^2)^{3/2}} \left( \frac{N}{l} \times (-R \operatorname{cosec}^2 \alpha d\alpha) \right) \vec{a}_z$$

$$= -\frac{IR^3 N}{2l} \operatorname{cosec}^2 \alpha d\alpha$$

$$\frac{1}{2l(R^2+z^2)^{3/2}}$$

$$= \frac{-IR^3 N}{2l} \int_{\alpha_1}^{\alpha_2} \frac{\operatorname{cosec}^2 \alpha d\alpha}{(R^2+z^2)^{3/2}} dz$$



$$\bar{H} = \frac{-\int R^3 N}{2l} \int_{\alpha_1}^{\alpha_2} \frac{\operatorname{cosec}^2 \alpha \, d\alpha}{(R^2 + R^2 \cot^2 \alpha)^{3/2}} \bar{a}_3$$

$$= \frac{-\int R^3 N}{2l R^3} \int_{\alpha_1}^{\alpha_2} \frac{\operatorname{cosec}^2 \alpha \, d\alpha}{(1 + \cot^2 \alpha)^{3/2}} \bar{a}_3$$

$$= \frac{-\int N}{2l} \int_{\alpha_1}^{\alpha_2} \frac{\operatorname{cosec}^2 \alpha}{(\operatorname{cosec}^2 \alpha)^{3/2}} \, d\alpha \bar{a}_3$$

$$= \frac{-\int N}{2l} \int_{\alpha_1}^{\alpha_2} \frac{1}{\operatorname{cosec} \alpha} \, d\alpha \bar{a}_3$$

$$\bar{H} = \frac{-\int N}{2l} \int_{\alpha_1}^{\alpha_2} \sin \alpha \, d\alpha \bar{a}_3$$

$$= \frac{-\int N}{2l} (-\cos \alpha)_{\alpha_1}^{\alpha_2} \bar{a}_3$$

$$\bar{H} = \frac{\int N}{2l} (\cos \alpha_2 - \cos \alpha_1) \bar{a}_3$$

case-2

$$\alpha_1 + \alpha_2 = \pi$$

$$\alpha_1 = \pi - \alpha_2$$

$$\cos(\alpha_1) = \cos(\pi - \alpha_2)$$

$$= -\cos \alpha_2$$

$$\bar{H} = \frac{\int N}{2l} (\cos \alpha_2 + \cos \alpha_2) \bar{a}_3$$

$$= \frac{\int N}{2l} 2 \cos \alpha_2 \bar{a}_3$$

$$\bar{H} = \frac{\int N}{l} \cos \alpha_2 \bar{a}_3$$

Case-2)-

when the length of the solenoid is

$$\alpha_2 = 0$$

$$\Rightarrow \cos \alpha_2 = \cos 0 \\ = 1$$

$$\vec{H} = \frac{IN}{l} \cos \alpha_2 \vec{a}_3$$

$$\vec{H} = \frac{IN}{l} \vec{a}_3$$

\* Relation b/n M.F.D (B), M.F.I (H) & Magnetic flux (φ):-

$$\text{B \& H} \\ \boxed{B = \mu H}$$

wkt

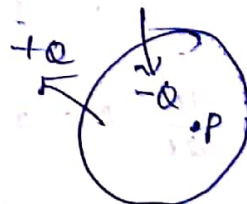
$$\vec{B} = \phi/A \text{ or } \phi/l_s \text{ or } \frac{\psi}{S}$$

$$\vec{B} = \frac{d\psi}{dS}$$

$$d\psi = \vec{B} \cdot d\vec{S}$$

$$\psi = \int_S \vec{B} \cdot d\vec{S}$$

\* Gauss law in Magnetostatics:-



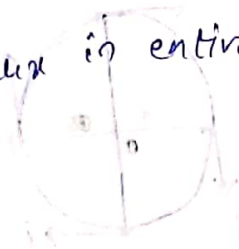
consider a closed surface in a magnetic field for every flux line enters into a surface. there must always be flux lines emerging "out of surface.



where total output flux = total input flux

- the outward flux is taken as + $\phi$  and inward as - $\phi$

then the magnetic flux in entire closed surface is equal to 0.



$$\therefore \Rightarrow \psi = 0$$

$$\boxed{\int_S \vec{B} \cdot d\vec{s} = 0}$$

\* Maxwell's second Equation ( $\nabla \cdot \vec{B} = 0$ ) :-

we know that the magnetic flux flowing through a surface is given by

$$\int \vec{B} \cdot d\vec{s} = \psi = \phi \rightarrow (1)$$

Acc. to Gauss law the Magnetic flux in entire closed surface is equal to zero.

$$\psi = 0 \rightarrow (2)$$

substitute (2) in (1)

$$\int_S \vec{B} \cdot d\vec{s} = 0 \rightarrow (3)$$

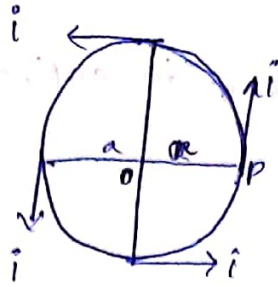
Apply the divergence theorem in (3)

$$\text{i.e., } \int_V \nabla \cdot \vec{B} \cdot d\vec{v} = 0$$

$$\nabla \cdot \vec{B} = 0$$

$$\boxed{\nabla \cdot \vec{B} = 0} \quad (or) \quad \boxed{(\text{div } \vec{B}) = 0}$$

\* Ampere's Circuit Law :-



statement:- The line integral of Magnetic field intensity ( $\vec{H}$ ) at any closed path [solenoid, square, conducting wire or circular coil] is equal to the current enclosed by that path.

$$\text{i.e., } \int_C \vec{H} \cdot d\vec{l} = I \quad \rightarrow (1)$$

Let current density  $J = I/s$

$$J = dI/ds$$

$$dI = J ds$$

$$I = \int_S J \cdot ds \quad \rightarrow (2)$$

substitute (2) in (1)

$$\int_C \vec{H} \cdot d\vec{l} = \int_S J \cdot ds = I$$

\* calculate the magnetic flux density at the centre of a current carrying loop when the radius of loop is 2cm, loop current is 1Amp and loop is placed in air

sol<sup>n</sup>: Given

$$I = 1 \times 10^3 \text{ Amp}$$

$$\text{radius of loop } a = 2 \text{ cm} = 2 \times 10^{-2} \text{ m}$$

$$\text{we know that } \vec{H} = \frac{I}{2a}$$

$$= \frac{10^3}{2 \times 2 \times 10^{-2}} = \frac{1}{40}$$

$$= 0.025 \text{ A-t/m}$$

$$\vec{B} = \mu_0 \vec{H} \quad [\text{air}]$$

$$= 4\pi \times 10^{-7} \times 0.025$$

$$= 0.314 \times 10^{-7}$$

$$= 3.14 \times 10^{-8}$$

=

$$\frac{\text{square}}{\pi d} = \frac{2I}{\pi d}$$

$$= \frac{\sqrt{2} \times 10^{-3}}{\pi \times 2 \times 10^{-2}}$$

$$= \frac{\sqrt{2}}{20 \times \pi}$$

$$= 0.022 \text{ A-t/m}$$

$$= 0.022 \text{ A-t/m}$$

\* the magnitude of  $\vec{H}$  at a radius of 1m from along linear conductor is 1A-t/m. find current in wire.

sol<sup>n</sup>: magnetic field intensity  $\vec{H} = \frac{I}{2a}$

$$\vec{H} = 1 \text{ A-t/m}, a = 1 \text{ m}$$

Ampere's circuit law

$$\int_S \vec{H} \cdot d\vec{l} = I$$

$$H \cdot \int dl = I \Rightarrow H (2\pi r) = I$$

$$I = 6.38 \text{ Amp}$$

\* The magnitude of  $H$  at a radius of  $4\text{ m}$  from a long linear conductor is  $2\text{ A-T/m}$ . find the current.

sol:-

$$r = 4\text{ m}$$

$$H = 2$$

$$\oint \vec{H} \cdot d\vec{l} = I$$

$$I = H \int dl$$

$$I = 2 \times 2\pi r \times 2$$

$$= 2 \times 2\pi \times 4 \times 2$$

$$= 50.26\text{ Amp}$$

\* A circular coil of radius  $1.5\text{ cm}$  carries a current  $1.5\text{ amps}$ . If the coil has  $25$  turns find the field at centre.

sol:-

$$H = \frac{NI}{2a}$$

$$= \frac{25 \times 1.5}{2 \times 1.5 \times 10^{-2}}$$

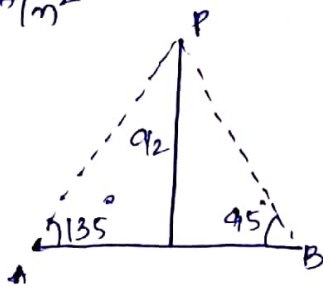
$$\vec{B} = \mu H = \frac{2}{4\pi \times 10^{-7}} \times 25 \times 1.5$$

$$= 157.07 \times 10^5 \text{ wb/m}^2$$

\* using Amphere's circuit law find  $\vec{H}$  and  $\vec{B}$  inside a long straight magnetic conductor of radius  $r$  mm

\* a) A steady current of  $I$  amp flow in a conductor bent in the form of square loop of amt. find magnetic field intensity at the centre of loop.

b) using Amphere's circuit law find  $\vec{H}$  and  $\vec{B}$  inside a long straight magnetic conductor of radius  $8$  mm carrying a current density of  $50 \text{ kA/m}^2$



Sol: i) 
$$\vec{H} = \frac{I}{4\pi d} \cos\theta \, a\vec{r}$$

$$= \frac{I}{4\pi d} (\cos\alpha_1 - \cos\alpha_2)$$

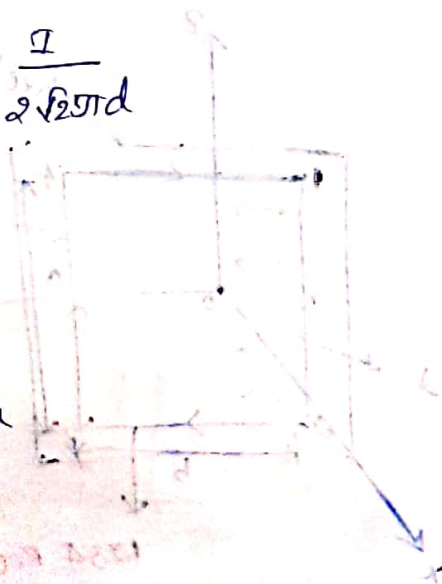
$$= \frac{I}{4\pi d} [\cos 45^\circ - \cos 135^\circ]$$

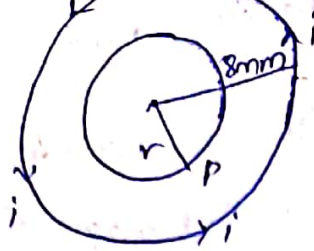
$$= \frac{I}{4\pi d} \left[ \frac{1}{\sqrt{2}} + \frac{1}{\sqrt{2}} \right]$$

$$= \frac{I}{4\pi d} \cdot \frac{2}{\sqrt{2}} = \frac{I}{2\sqrt{2}\pi d}$$

ii) 
$$\vec{H} = 4\vec{H} = 4 \cdot \frac{I}{2\sqrt{2}\pi a}$$

$$= \frac{\sqrt{2}I}{\pi a} \Rightarrow \vec{H} \times$$





$$\int \vec{H} \cdot d\vec{l} = i$$

$$\vec{H} \cdot dl = J \cdot A$$

$$H \cdot \int dl = J \cdot A$$

$$H \cdot (2\pi r) = J \cdot \pi r^2$$

$$H \cdot 2 = J \cdot r$$

$$= \frac{50 \times 10^3 \times 4 \times 10^{-3}}{2}$$

$$H = 200 \text{ A/m}$$

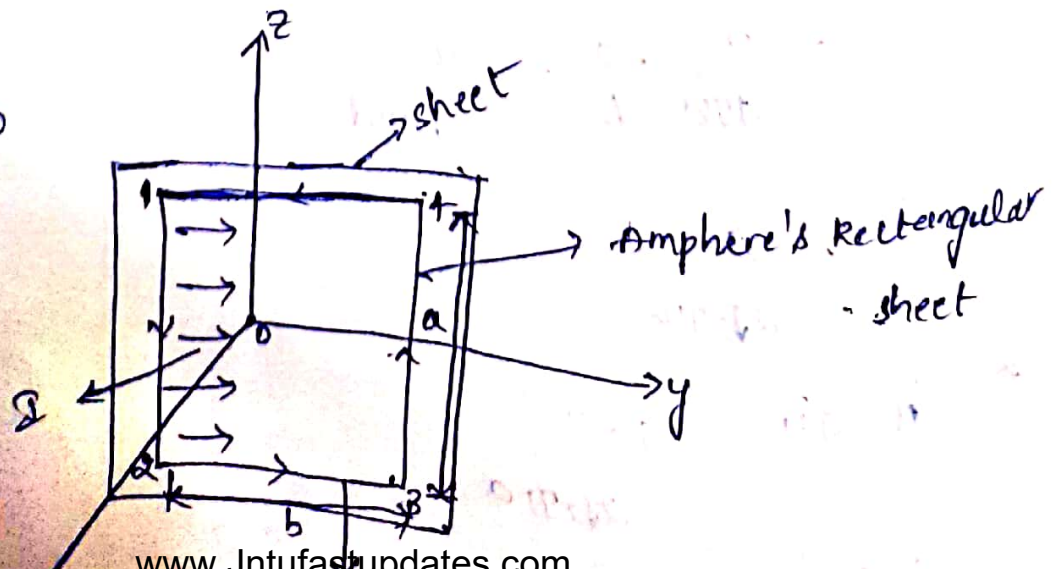
$$B = \mu \cdot H$$

$$= 4\pi \times 10^{-7} \times 200$$

$$= 2.51 \times 10^{-4} \text{ Tesla}$$

MFI due to infinite sheet of charge :-

xy plane  
z=0



- Here we consider the infinite sheet in  $z=0$  plane nothing but  $x-y$  plane sheet.
- Take the Amphere's surface. if the sheet place that at the top and bottom.
- Amphere's surface sheet or 1234 sheet i.e., is rectangular form. by using this sheet we are cutting that infinite sheet ( $x-y$  - plane sheet). that time the current direction will be shown in figure. i.e., nothing but current density.
- we know that the Amphere's Circuit Law

$$\oint \vec{H} \cdot d\vec{l} = I \rightarrow (1)$$

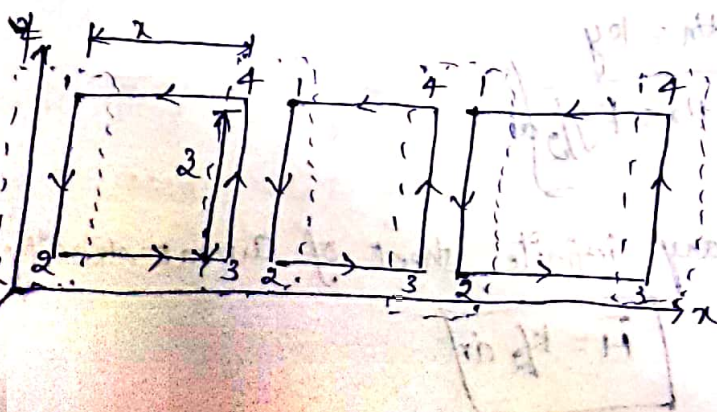
Applying Amphere's circuit law to the rectangular sheet is given by,

$$\oint \vec{H} \cdot d\vec{l} = ky \cdot b = I \rightarrow (2)$$

where  $ky =$  Surface Current density

$b =$  length

1234 - Infinite sheets



→ 3-4 sides in Amphere's surface is in up direction of vector and 1-2 Amphere's surface sheet is in down ward direction then both are in opposite in direction.

$$\int \vec{H} \cdot d\vec{l} = 0$$

∴ The total magnetic field intensity in rectangular sheet (1234) is,

$$\int_1^2 \vec{H} \cdot d\vec{l} + \int_2^3 \vec{H} \cdot d\vec{l} + \int_3^4 \vec{H} \cdot d\vec{l} + \int_4^1 \vec{H} \cdot d\vec{l} = ky \cdot b$$

$$\int_2^3 \vec{H} \cdot d\vec{l} + \int_4^1 \vec{H} \cdot d\vec{l} = ky \cdot b$$

$$\int_2^3 H_x dx + \int_4^1 H_x dx = ky \cdot b$$

$$H_x \int_2^3 dx + H_x \int_4^1 dx = ky \cdot b$$

the integration of dx is b in figure. [∫ dx = b]

$$\therefore H_x b + H_x b = ky \cdot b$$

$$2H_x b = ky \cdot b$$

$$2H_x = ky$$

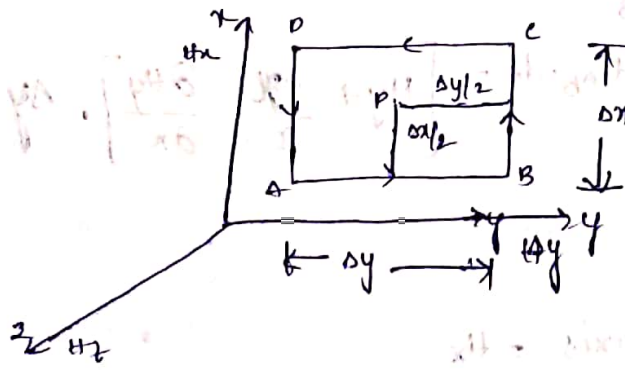
$$H_x = \frac{ky}{2} \hat{a}_x$$

In general for any infinite sheet of current density is

$$\vec{H} = \frac{k}{2} \vec{a}_x$$



# Point form of Ampere's Circuit law:-



- consider a point P the magnetic field intensity be  $H$  and its components be  $H_x$ ,  $H_y$  and  $H_z$ .
- The values of both  $H$  and its components change with distance as moving move in away direction.
- consider a rectangular loop ABCD parallel to xy-plane. such that point-P is located exactly its centre, so in order to evaluate the closed line integral of ' $H$ ' over the path in the direction of ABCD, it can be divided into 4 parts namely AB, BC, CD & DA segments/sides.

AB segment:-

with the help of Taylor's Series theorem.

$$AB \parallel \text{to } y\text{-axis} = H_y$$

$$\text{at distance} = \Delta y$$

$$\text{at point P} = \Delta x/2$$

The magnetic field intensity AB segment,

$$H_{AB} = H_y + \frac{\Delta x}{2} \frac{\partial H_y}{\partial x}$$

the line integral of AB is,

$$\int_A^B H_{AB} \cdot dl = \left[ Hy + \frac{\Delta y}{2} \frac{\partial Hy}{\partial x} \right] \cdot \Delta y \rightarrow (2)$$

BC segment:-

$$BC \parallel \text{to } x\text{-axis} = Hx$$

$$\text{Distance} = \Delta x$$

$$\text{at point } P = \frac{\Delta y}{2}$$

with by using Taylor's Series theorem

$$H_{BC} = Hx + \frac{\Delta y}{2} \frac{\partial Hx}{\partial y}$$

$$\int_B H_{BC} \cdot dl = \left[ Hx + \frac{\Delta y}{2} \frac{\partial Hx}{\partial y} \right] \Delta x \rightarrow (2)$$

CD segment:-

$$CD \parallel y\text{-axis} = Hy$$

$$\text{Distance} = \Delta y$$

$$\text{at point } P = \Delta x/2$$

The magnetic field intensity of CD segment is,

$$H_{CD} = -Hy + \frac{\Delta x}{2} \frac{\partial Hy}{\partial x}$$

$$\int_C^D H_{CD} \cdot dl = \left[ -Hy + \frac{\Delta x}{2} \frac{\partial Hy}{\partial x} \right] \cdot \Delta y \rightarrow (3)$$

DA segment :-

DA || to x-axis =  $H_x$

Distance =  $\Delta x$

at point P =  $-dy/2$

By using Taylor's series theorem

The magnetic field intensity of DA segment is

$$H_{DA} = H_x - \frac{\Delta y}{2} \frac{\partial H_x}{\partial y}$$

$$\int_D^A H_{DA} \cdot dl = \left[ H_x - \frac{\Delta y}{2} \frac{\partial H_x}{\partial y} \right] \cdot \Delta x \rightarrow (4)$$

The line integral of ABCD rectan<sup>g</sup>le loops MFI ( $H$ ) of ABCD rectangle

$$\text{Loop's } \int H_{ABCD} \cdot dl = \int_A^B H_{AB} \cdot dl + \int_B^C H_{BC} \cdot dl + \int_C^D H_{CD} \cdot dl + \int_D^A H_{DA} \cdot dl$$

$$= H_y \Delta y + \frac{\Delta x}{2} \frac{\partial H_y}{\partial x} \Delta y - H_x \Delta x - \frac{\Delta y}{2} \frac{\partial H_x}{\partial y} \Delta x - H_y \Delta y + \dots$$

$$\frac{\Delta x}{2} \frac{\partial H_y}{\partial x} \Delta y + H_x \Delta x - \frac{\Delta y}{2} \frac{\partial H_x}{\partial y} \Delta x$$

$$= \Delta x \Delta y \frac{\partial H_y}{\partial x} - \Delta x \Delta y \frac{\partial H_x}{\partial y}$$

$$H_{ABCD} \text{ all} = \Delta x \Delta y \left[ \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} \right]$$

$$\boxed{I = \Delta x \Delta y \left[ \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} \right]}$$

current density  $J = \delta/A$

$$I = J \cdot A$$

$$J \cdot A = \delta x \cdot \delta y \left[ \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} \right]$$

$$\text{let } A = \delta x \cdot \delta y$$

$$J \cdot [\delta x \cdot \delta y] = \delta x \cdot \delta y \left[ \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} \right]$$

$$J = \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} \rightarrow (5)$$

$$\boxed{J = \nabla \times \vec{H}}$$

or

Now applying the unit vector for each component

$[\hat{a}_x, \hat{a}_y, \hat{a}_z]$

$$\Rightarrow \left[ \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} \right] \hat{a}_z + \left[ \frac{\partial H_z}{\partial y} - \frac{\partial H_y}{\partial z} \right] \hat{a}_x + \left[ \frac{\partial H_x}{\partial z} - \frac{\partial H_z}{\partial x} \right] \hat{a}_y = J$$

It can be expressed in matrix form as

$$\begin{vmatrix} \hat{a}_x & \hat{a}_y & \hat{a}_z \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ H_x & H_y & H_z \end{vmatrix} = J$$

$$\therefore \boxed{\nabla \times \vec{H} = J}$$

The above equation is also called as Maxwell's third Equation.

∴ Maxwell's Third Equation  $[\nabla \times \vec{H} = \vec{J}]$  :-

We know that

According to Ampere's circuit law

$$\int_C \vec{H} \cdot d\vec{l} = I \quad \rightarrow (1)$$

Let current density  $\vec{J} = I/s$

$$I = \vec{J} \cdot s$$

$$dI = \vec{J} \cdot d\vec{s}$$

$$I = \int_S \vec{J} \cdot d\vec{s} \quad \rightarrow (2)$$

Substitute (2) in (1)

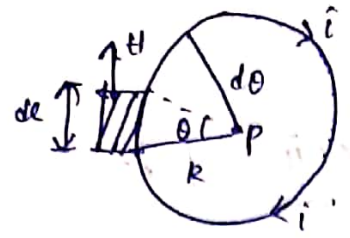
$$\int_C \vec{H} \cdot d\vec{l} = \int_S \vec{J} \cdot d\vec{s}$$

Now applying Stokes theorem

$$\int_C \nabla \times \vec{H} \cdot d\vec{s} = \int_S \vec{J} \cdot d\vec{s}$$

$$\boxed{\nabla \times \vec{H} = \vec{J}}$$

\* Magnetic field due to Circular loop:-



from the figure we see

$$dl = R d\theta$$

from Biot-Savart's law

$$d\vec{H} = \frac{I dl \sin\theta}{4\pi R^2} \vec{ar}$$

$$= \frac{I R d\theta \sin\theta}{4\pi R^2} \vec{ar}$$

If  $\theta = 90^\circ$

$$\sin\theta = 1$$

$$d\vec{H} = \frac{I d\theta}{4\pi R} \vec{ar}$$

$$H = \int_0^{2\pi} \frac{I d\theta}{4\pi R} \vec{ar}$$

$$= \frac{I \vec{ar}}{4\pi R} \int_0^{2\pi} d\theta$$

$$= \frac{I \vec{ar}}{4\pi R} [2\pi - 0]$$

$$\vec{H} = \frac{I \vec{ar}}{2R} \text{ Amp/m}$$

Magnetic flux density

$$\vec{B} = \mu H$$

$$\vec{B} = \frac{\mu I}{2R} \vec{a}_r \quad \text{Tesla}$$

# Force in Magnetic fields

Introduction :-

Magnetic field :-

The study of interaction between the moving charges in moving position is called as Magnetic field.

The Electric charges at rest produce electric field, the charges at moving produce magnetic field.

However, magnetic field produce a moving charges only when charged particles having charge  $Q$  travels with velocity  $v$  in a magnetic field  $B$  it experience a force that force is called as Magnetic force [ $f_m$ ].

$$\therefore \boxed{f_m = Q(\vec{v} \times \vec{B})}$$

The above equation is called as Magnetic force equation. if  $\vec{v}$  Parallel to  $\vec{B}$  then force is 0.

$$\vec{v} \parallel \vec{B} \Rightarrow \boxed{f_m = 0}$$

By Taking The Magnitudes in equation the force is proportional to  $v$  and  $B$ .

Similarly The force is proportional to  $\sin \theta$ , because we take the magnitudes.

$$\boxed{f_m = BvQ \sin \theta}$$



## \* force on a moving charge / Lorentz force Equation:-

In electrostatic field ( $\vec{E}$ ) The force on a charged particle is,

$$f_e = qE \rightarrow (1)$$

A charged particle in motion in a magnetic field of flux density  $\vec{B}$  to experience a force is given by,

$$f_m = q(\vec{v} \times \vec{B}) \rightarrow (2)$$

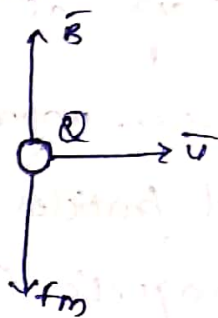
There is a fundamental difference between electric field and Magnetic field.

$\therefore$  The force is given by,

$$f = f_e + f_m$$

$$= q\vec{E} + q(\vec{v} \times \vec{B})$$

$$f = q[\vec{E} + (\vec{v} \times \vec{B})]$$



where  $v$  = velocity of charged particles

$\vec{E}$  = Electric field intensity

$\vec{B}$  = Magnetic flux density

$q$  = charge

$\therefore$  The above equation is called as force on a moving charge / Lorentz force Equation.

\* force on a current element in Magnetic field:-

We know that force on a Magnetic field.

$$f_m = Q(\vec{v} \times \vec{B})$$

Now we consider line charge, surface charge and volume charge.

then force on a current element in a M.F is,

$$\vec{J} = \rho_v \cdot \vec{v} \longrightarrow (1)$$

$$\vec{J} \cdot d\vec{u} = k \cdot ds = \vec{J} \cdot d\vec{l} \longrightarrow (2)$$

$\vec{J}$  = current density

$k$  = surface charge density [C/m<sup>2</sup>]

$\vec{J} \cdot d\vec{l}$  = current element

let

$$\vec{J} \cdot d\vec{u} = \vec{J} \cdot d\vec{l} \longrightarrow (3)$$

substitute (1) in (3)

$$\rho_v \cdot \vec{v} \cdot d\vec{u} = \vec{J} \cdot d\vec{l}$$

$$\rho_v d\vec{u} \cdot \vec{v} = \vec{J} \cdot d\vec{l}$$

$$\boxed{dQ \cdot \vec{v} = \vec{J} \cdot d\vec{l}} \longrightarrow (4)$$

$$\rho_v = Q/v = \frac{dQ}{dv}$$

$$dQ = \rho_v \cdot dv$$

$$\text{hence } f_m = Q \cdot \vec{v} \times \vec{B} \longrightarrow$$

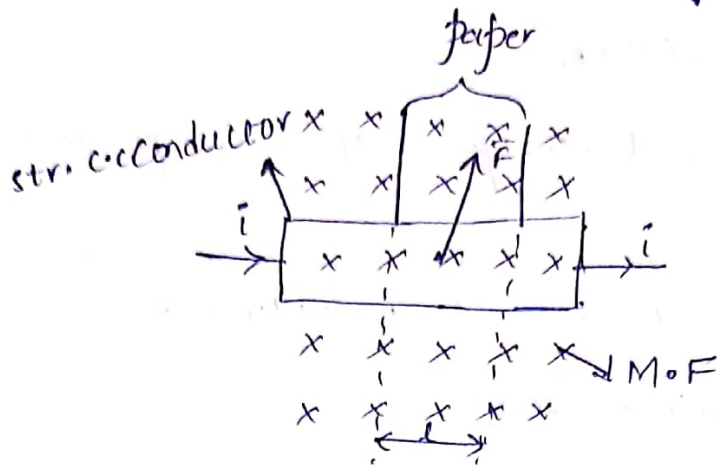
$$df_m = dQ \cdot \vec{v} \times \vec{B} \longrightarrow (5)$$

substitute (4) in (5)

$$df_m = dQ \vec{J} \cdot d\vec{l} \times \vec{B} \quad \text{or} \quad df_m = k ds \times \vec{B} \quad df_m = \vec{J} \cdot d\vec{l} \times \vec{B}$$

$$f_m = \int \vec{J} \cdot d\vec{l} \times \vec{B} \quad f_m = \int k \cdot ds \times \vec{B} \quad f_m = \int \vec{J} \cdot d\vec{l} \times \vec{B}$$

Force on a straight and long current carrying conductor in a magnetic field (or) filamentary current.



WKT force on M.F is

$$F_m = q(\vec{v} \times \vec{B}) \rightarrow (1)$$

WKT  $\vec{v} = \frac{dl}{dt} \rightarrow (2)$

substitute (2) in (1)

$$F_m = q \frac{dl}{dt} \times \vec{B}$$

$$dF_m = dq \frac{dl}{dt} \times \vec{B}$$

$$dF_m = \frac{dq}{dt} dl \times \vec{B}$$

$$dF_m = I dl \times \vec{B}$$

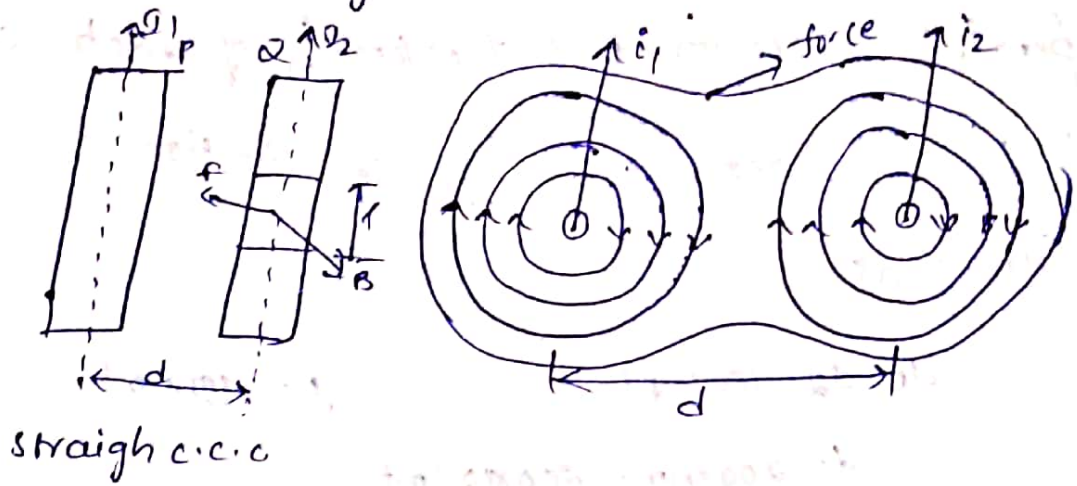
$$F_m = BIL \sin \theta$$

where  $\theta$  represents the angle b/w vectors.

i.e., direction of current flow and direction of  $\vec{B}$ .

$\therefore$  The above equation is called as filamentary/familiar

\* force on a 2-straight and parallel current carrying conductors due to magnetic field.



WKT Magnetic field intensity due straight c.c.c is,

$$\vec{H} = \frac{I_1}{2\pi d} \quad [\text{for p-conductor}]$$

WKT  $\vec{B} = \mu \vec{H}$

$$\vec{B} = \frac{\mu I_1}{2\pi d} \rightarrow W$$

force on conductor Q, due to P is

$$f = B L I_2 \rightarrow (2)$$

$$f = \frac{\mu I_1 L I_2}{2\pi d}$$

$$\boxed{\frac{f}{L} = \frac{\mu I_1 I_2}{2\pi d} \quad N/m}$$

→ 2 long parallel conduction carrying 100amps. If the conductors are separated by 200mm. find the force/m of each conductor if the current flow direction is in opposite direction.

sol: Given that

$$I_1 = I_2 = 100 \text{ amps}$$

$$\mu = 4\pi \times 10^{-7}$$

$$d = 200 \text{ mm} = 200 \times 10^{-3} \text{ m}$$

$$F = ?$$

$$\frac{F}{L} = \frac{\mu I_1 I_2}{2\pi d} = \frac{4\pi \times 10^{-7} \times 100 \times 100}{2\pi \times 0.2}$$

$$= 0.1 \text{ N/m}$$

→ A Galvanometer has a rectangular coil suspended in a radial magnetic field so that the magnetic field always act across the plane of coil of 10mm by 10mm side and has 1000 turns if the magnetic provides a magnetic flux density of 0.3 T. find the torque entered on the coil for a current of 10mA.

sol: Given that

$$A = 10 \times 10 \text{ mm}$$

$$N = 1000$$

$$= 100 \times 10^{-6}$$

$$T = ?$$

$$B = 0.3 \text{ T}$$

$$I = 10 \text{ mA} =$$

$$\text{torque} = \underline{M \cdot B} \text{ magnetic dipole moment}$$

$$M = I \cdot A = N I A = 1000 \times 10 \times 10^{-6} \times 10 \times 10^{-6}$$

$$T = 0.3 \times 10^{-3}$$

$$= 3 \times 10^{-4} \text{ N/m}$$

\* A distributional line consist of 2 straight parallel conductors supported on the cross arms of wooden poles spaced 100m. the normal spacing b/w the two conductors is 20cm suppose a current of 10,000 amp. determine the force/mt.

sol<sup>n</sup> Given that  $I_1 I_2 = 10,000$   
 $d = 20 \text{ cm} = 20 \times 10^{-2} \text{ m}$   
 $l = 100 \text{ m}$

$$F = \frac{\mu I_1 I_2 l}{2\pi d} = \frac{4\pi \times 10^{-7} \times 10,000 \times 10,000 \times 100}{2 \times \pi \times 20 \times 10^{-2}}$$

$$= 10^4 \text{ N.}$$

\* The force b/w 2 long parallel conductors is 15 N/m. The conductor spacing of 10cm. if one conductor carries twice the current of other: calculate the current in each conductor.

sol<sup>n</sup>  $f/l = 15 \text{ N/m}$

$$d = 10 \text{ cm} = 10 \times 10^{-2} \text{ m}$$

$$I_2 = 2I_1$$

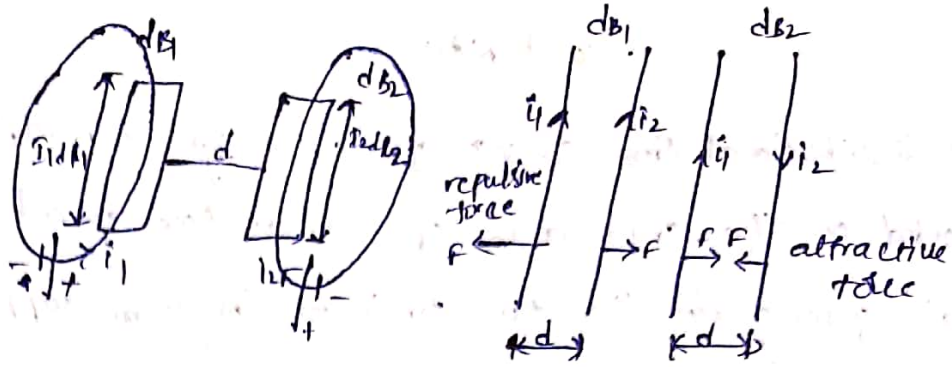
$$\frac{f}{l} = \frac{\mu \cdot I_1 \cdot I_2}{2\pi d}$$

$$15 = \frac{4\pi \times 10^{-7} \times I_1 \times 2I_1^2}{2 \times \pi \times 10^{-1}} = 40 \times 10^{-7} \times I_1^2$$

$$I_1^2 = \frac{15}{40 \times 10^{-7}} = \frac{10^6 \times 15}{40} = 3.75 \times 10^4$$

$$I_1 = 1936 \quad I_2 = 3872$$

\* force b/w Two differential Current Elements:-



→ let us consider two current elements  $i_1 dl_1$ ,  $i_2 dl_2$  as shown in-figure.

- Note that the directions of  $i_1$  &  $i_2$  are same. the current elements produced magnetic fields ( $dB_1$  &  $dB_2$ )

- As the currents are flowing in same direction through the elements the force will be developed ( $d(dF_1)$ ) on element  $i_1 dl_1$  due to magnetic field  $dB_2$ .

- As the currents are flowing in opposite direction through the elements the force will be developed ( $d(dF_2)$ ) on element  $i_2 dl_2$  due to magnetic field  $dB_1$ .

from the equations of the force on a differential current element is  $d(dF_1) = i_1 dl_1 \times dB_2 \rightarrow (a)$

Acc. to Biot-Savart's law

$$dB_2 = \frac{\mu_0 i_2 dl_2}{4\pi r_2^2} \cdot \hat{a}_{r_2} \rightarrow (1)$$

W.K.T.

$$\bar{B} = \mu \bar{H}$$

$$d\bar{B}_2 = \mu d\bar{H}_2 \rightarrow (2)$$

substitute (1) in (2)

$$d\bar{B}_2 = \frac{\mu I_2 dl_2}{4\pi R_2^2} \cdot \bar{a}_{r_{21}} \rightarrow (3)$$

substitute (3) in (a)

$$d(d\bar{f}_1) = I_1 dl_1 \times \frac{\mu I_2 dl_2}{4\pi R_2^2} \bar{a}_{r_{21}}$$

$$d\bar{f}_1 = \int_{l_1} \frac{\mu I_1 I_2 dl_1 d\bar{b}_2}{4\pi R_2^2} \bar{a}_{r_{21}}$$

$$f_1 = \iint_{l_1 l_2} \frac{\mu I_1 I_2 \cdot dl_1 d\bar{b}_2}{4\pi R_2^2} \bar{a}_{r_{21}}$$

Exactly, following the same steps we can conclude that,

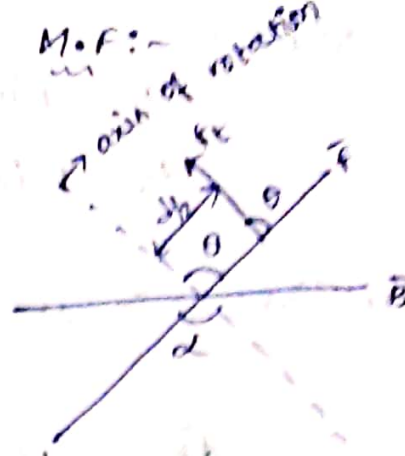
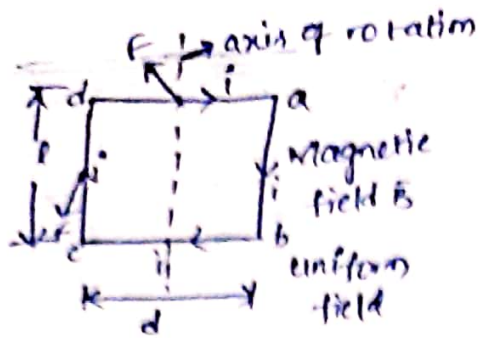
$$d(d\bar{f}_2) = I_2 d\bar{b}_2 \times d\bar{B}_1$$

$$f_2 = -f_1$$

Thus the above equation indicates that both the forces  $f_1$  &  $f_2$  obey Newton's 3<sup>rd</sup> law.



\* Torque on a current loop in a M.F.:-



AKT

The force on a current element in a magnetic field

$$d\vec{F} = I(d\vec{l} \times \vec{B})$$

$$\vec{F} = BIL \sin \theta$$

consider a rectangular loop ABCD with a sides of  $L$  &  $d$  carrying current of  $i$  amps is placed in a uniform magnetic field ( $\vec{B}$ ). i.e., shown in figure.

The force on each element of a loop side 'AB' of the length

$l$  m

$$i.e., d\vec{F} = I(d\vec{l} \times \vec{B})$$

$$\vec{F} = BIL \sin \theta$$

$$\theta = 90^\circ$$

$$\vec{F} = BIL \rightarrow (1)$$

\* If the plane of the loop is made an angle  $\theta$  i.e., shown in figure. the tangential force is

$$F_t = F \cos \theta$$

$$F_t = BIL \cos \theta \rightarrow (2)$$

the total torque on the loop is

$$\tau = 2x (\text{torque on each side})$$

$$= 2x (f \ell \times d/2)$$

$$= f \ell \times d$$

$$= B \ell \cos \theta \times d$$

$$= B \ell \cos \theta (\ell \times d)$$

$$= B \ell A \cos \theta$$

$$= B(\ell A) \cos \theta$$

$$= B m \cos \theta$$

$$\boxed{\tau = B m \cos \theta}$$

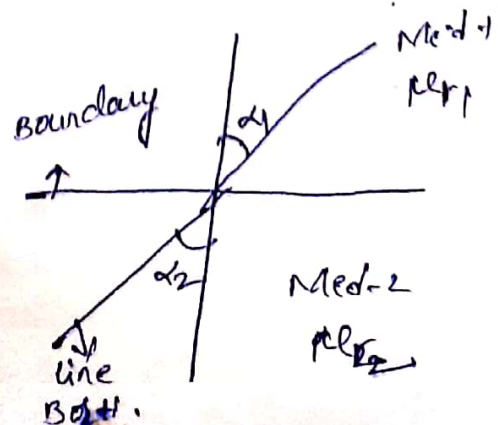
If  $\theta = 0$

$$\boxed{\tau = B m}$$

$\vec{m}$  = dipole moment

\* Derive the boundary conditions at the magnetic interfaces and show that  $\frac{\tan \alpha_1}{\tan \alpha_2} = \frac{\mu_1}{\mu_2}$ . The figure shows a magnetic field line passing through the interface between two magnetic media of permeabilities  $\mu_1$  and  $\mu_2$ .

in order to develop the boundary at an angle  $\alpha_1$  in media-1.  $\alpha_2$  is media-2. let us assume that both media are having  $\mu$  and  $B$  are in same direction and  $\vec{B} = \mu \vec{H}$  is valid



<sup>n</sup>  
sol: Now the relations for magnetic boundary,

$$B_1 \cos \alpha_1 = B_2 \cos \alpha_2 \rightarrow (1)$$

$$\mu_1 \sin \alpha_1 = \mu_2 \sin \alpha_2 \rightarrow (2)$$

$$(1) \Rightarrow B_1 \cos \alpha_1 = B_2 \cos \alpha_2$$

$$\mu_0 \mu_1 H_1 \cos \alpha_1 = \mu_0 \mu_2 H_2 \cos \alpha_2 \rightarrow (3)$$

$$(2)/(3) \Rightarrow \frac{\mu_1 \sin \alpha_1}{\mu_0 \mu_1 H_1 \cos \alpha_1} = \frac{\mu_2 \sin \alpha_2}{\mu_0 \mu_2 H_2 \cos \alpha_2}$$

$$\frac{\tan \alpha_1}{\mu_1} = \frac{\tan \alpha_2}{\mu_2} \quad \left[ \begin{array}{l} \mu_1 \rightarrow \mu \\ \mu_2 \rightarrow \mu \end{array} \right]$$

$$\boxed{\frac{\tan \alpha_1}{\tan \alpha_2} = \frac{\mu_1}{\mu_2}}$$

$$\boxed{\frac{\tan \theta_1}{\tan \theta_2} = \frac{\mu_1}{\mu_2}}$$